

# The Sivers effect in $p\bar{p}$ -interactions

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- Introduction
- Parton densities (forward,  $p_T$ -dependent)
- Semi-inclusive processes
- Sivers effect and single-spin asymmetries
- Phenomenology of the Sivers effect
- Drell-Yan measurements in  $p\bar{p}$ -interactions
- Summary

# General strategy

## 1. Factorization

$$\sigma \propto (\text{pert. part}) \otimes (\text{non-pert. part}) + \mathcal{O}\left(\frac{1}{Q}\right)$$

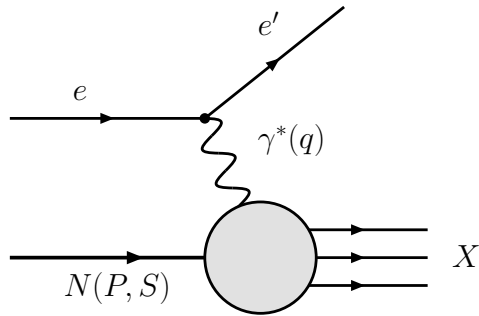
- required for interpretation
- does not hold for each hard process
- required scale:  $Q > 1 \text{ GeV}$

## 2. Universality

- allows one to make predictions
- can be non-trivial

# Inclusive deep-inelastic scattering

- Process



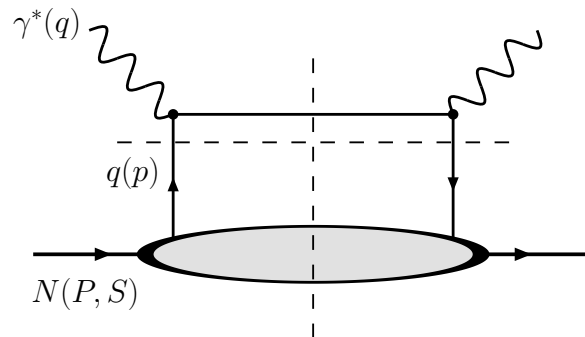
- Variables

$$Q^2 = -q^2 > 0 \quad x = \frac{Q^2}{2P \cdot q} \quad \Phi_l^S$$

- Optical theorem

$$\sigma_{\gamma^* N \rightarrow X} \propto \text{Im} A(\gamma^* N \rightarrow \gamma^* N, \vartheta = 0)$$

- Parton model



- Factorization

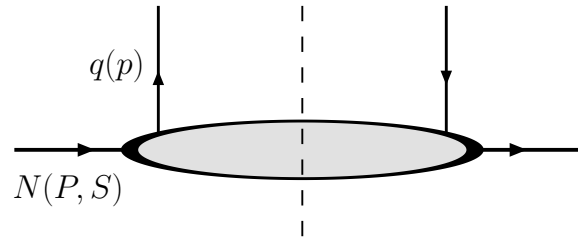
e.g.,  $f_1^q(x, Q^2)$

$$p^+ = \frac{1}{\sqrt{2}}(p^0 + p^3) = xP^+$$

$$p^- = \frac{1}{\sqrt{2}}(p^0 - p^3), \quad \vec{p}_T = (p^1, p^2)$$

# Parton densities

## 1. Forward parton densities



- Unpolarized pdf

$$f_1(x) = \int \frac{d\xi^-}{4\pi} e^{ixP^+\xi^-} \langle P, S | \bar{\psi}(0) \gamma^+ \psi(\xi^-) | P, S \rangle$$

$$\xi^- = \frac{1}{\sqrt{2}}(t - z) \quad \xi^+ = \xi_T = 0$$

- Helicity pdf

$$g_1(x) \propto \langle | \bar{\psi}(0) \gamma^+ \gamma_5 \psi(\xi^-) | \rangle$$

- Transversity pdf

$$h_1(x) \propto \langle | \bar{\psi}(0) i\sigma^{i+} \gamma_5 \psi(\xi^-) | \rangle \quad \sigma^{\mu\nu} = \frac{i}{2}[\gamma^\mu, \gamma^\nu]$$

2.  $p_T$ -dependent parton densities  
(Boer, Mulders, 1998)

$$\int \frac{d\xi^- d^2\vec{\xi}_T}{2(2\pi)^3} e^{i(xP^+\xi^- - \vec{p}_T \cdot \vec{\xi}_T)} \langle P, S | \bar{\psi}(0) \gamma^+ \psi(\xi^-, \vec{\xi}_T) | P, S \rangle$$

$$= f_1(x, \vec{p}_T^2) - \frac{\varepsilon_T^{ij} p_{Ti} S_{Tj}}{M} f_{1T}^\perp(x, \vec{p}_T^2)$$

Parameterization of

$$\Phi^{[\Gamma]}(x, \vec{p}_T, S) = \int \frac{d\xi^- d^2\vec{\xi}_T}{2(2\pi)^3} e^{i(xP^+\xi^- - \vec{p}_T \cdot \vec{\xi}_T)} \langle P, S | \bar{\psi}(0) \Gamma \psi(\xi^-, \vec{\xi}_T) | P, S \rangle$$

Leading order (8 functions, 2 T-odd)

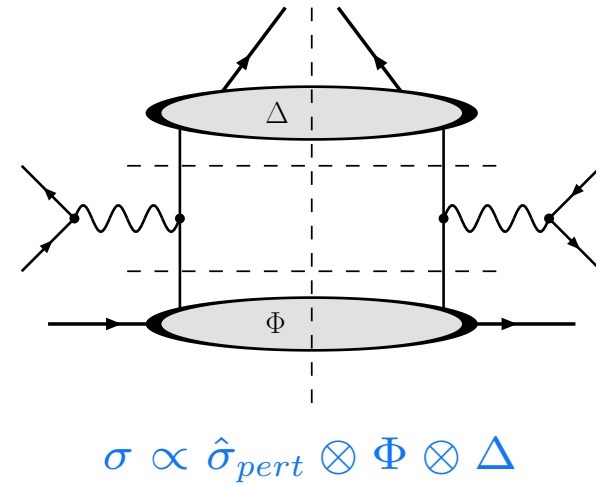
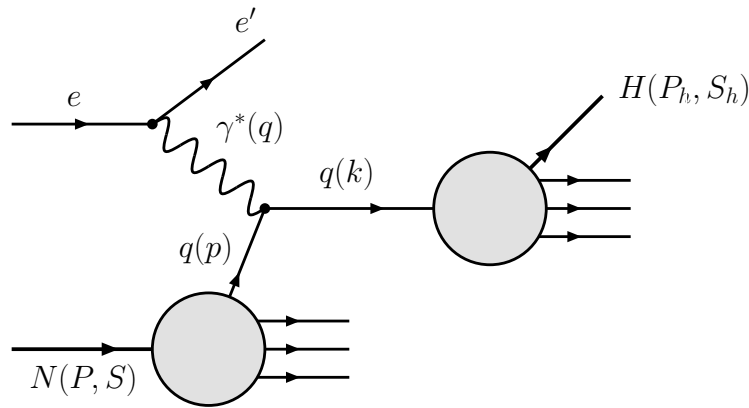
$$\Phi^{[\gamma^+]} = f_1 - \frac{\varepsilon_T^{ij} p_{Ti} S_{Tj}}{M} f_{1T}^\perp$$

$$\Phi^{[\gamma^+ \gamma_5]} = \lambda g_{1L} + \frac{\vec{p}_T \cdot \vec{S}_T}{M} g_{1T}$$

$$\Phi^{[i\sigma^{i+} \gamma_5]} = S_T^i h_{1T} + \frac{p_T^i}{M} \left( \lambda h_{1L}^\perp + \frac{\vec{p}_T \cdot \vec{S}_T}{M} h_{1T}^\perp \right) - \frac{\varepsilon_T^{ij} p_{Tj}}{M} h_1^\perp$$

# Semi-inclusive processes

## 1. Semi-inclusive deep-inelastic scattering $e N \rightarrow e' H X$



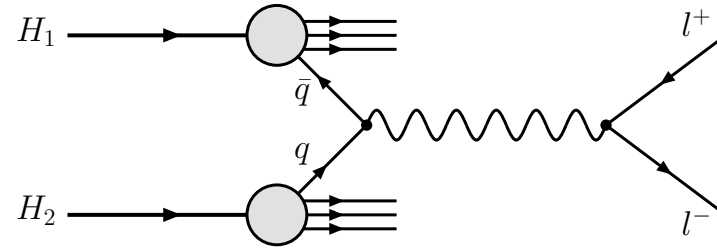
Factorization formula depends on kinematics

- $\sigma$  integrated upon  $P_{hT}$
- $\sigma$  differential in  $P_{hT}$  and  $P_{hT} \approx Q$
- $\sigma$  differential in  $P_{hT}$  and  $P_{hT} \ll Q$

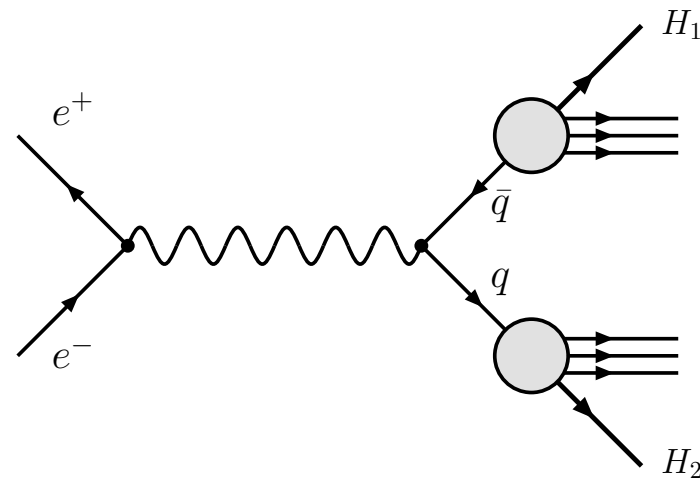
$$\sigma \propto \hat{\sigma}_{pert} \otimes \text{pdf} \otimes \text{ff} \otimes \text{soft}$$

(Collins, Soper, 1981; Collins, Soper, Sterman, 1985;  
Ji, Ma, Yuan, 2004; Collins, Metz, 2004)

2. Drell-Yan process  $H_1 H_2 \rightarrow l^+ l^- X$



3. Electron-positron annihilation  $e^+ e^- \rightarrow H_1 H_2 X$



## T-odd parton densities

- (Sivers, 1990)

$$\int \frac{d\xi^- d^2\vec{\xi}_T}{2(2\pi)^3} e^{i(xP^+\xi^- - \vec{p}_T \cdot \vec{\xi}_T)} \langle P, \uparrow | \bar{\psi}(0) \gamma^+ \psi(\xi^-, \vec{\xi}_T) | P, \uparrow \rangle$$
$$= f_1(x, \vec{p}_T^2) - \frac{(\vec{p}_T \times \vec{S}_T) \cdot \hat{P}}{M} f_{1T}^\perp(x, \vec{p}_T^2)$$

→ single-spin/azimuthal asymmetry

- (Boer, Mulders, 1998)

$$h_1^\perp(x, \vec{p}_T^2) \propto \langle P, \text{unp} | \bar{\psi}(0) i\sigma^{i+} \gamma_5 \psi(\xi^-, \vec{\xi}_T) | P, \text{unp} \rangle$$

Time-reversal:

$$f_{1T}^\perp = -f_{1T}^\perp = 0 \quad h_1^\perp = 0$$



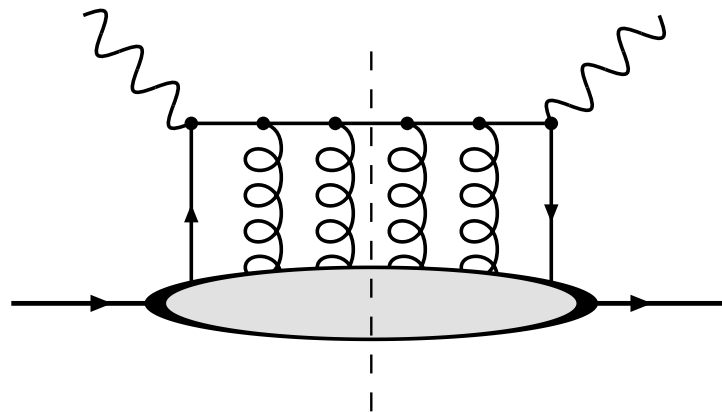
# Color gauge invariance

## 1. Forward partonic functions

$$\int d\xi^- e^{ixP^+\xi^-} \langle | \bar{\psi}(0) \Gamma \mathcal{L}(0; \xi^-) \psi(\xi^-) | \rangle$$

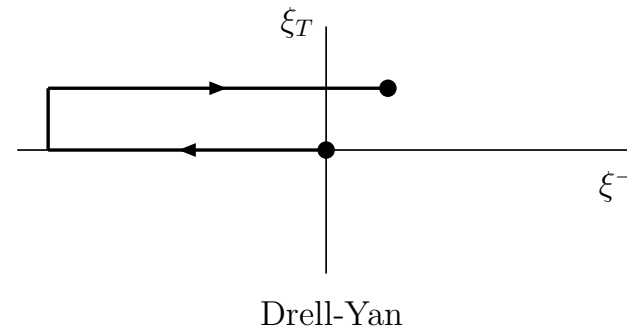
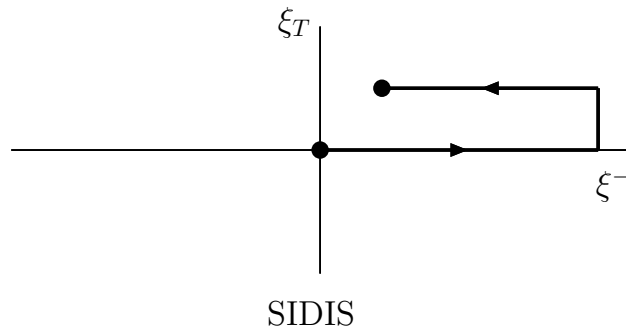
$$\mathcal{L}(0; \xi^-) = \mathcal{P} \exp \left( -ig \int_0^{\xi^-} d\eta^- A^+(0, \eta^-, \vec{0}_T) \right)$$

Gauge-link generated by rescattering



## 2. $p_T$ -dependent partonic functions

$$\int d\xi^- d^2\vec{\xi}_T e^{i(xP^+\xi^- - \vec{p}_T \cdot \vec{\xi}_T)} \langle | \bar{\psi}(0) \Gamma \mathcal{L}(0, \vec{0}_T; \xi^-, \vec{\xi}_T) \psi(\xi^-, \vec{\xi}_T) | \rangle$$



- $\mathcal{L}(0, \vec{0}_T; \xi^-, \vec{\xi}_T) = \mathcal{L}(0, \vec{0}_T; \infty, \vec{0}_T) \times \mathcal{L}(\infty, \vec{0}_T; \infty, \vec{\xi}_T) \times \mathcal{L}(\infty, \vec{\xi}_T; \xi^-, \vec{\xi}_T)$   
(Belitsky, Ji, Yuan, 2002)
- Gauge-link  $\rightarrow$  T-odd pdfs can be non-zero  
(Brodsky, Hwang, Schmidt, 2002; Collins, 2002)
- Different links for semi-inclusive DIS and Drell-Yan  $\rightarrow$  Universality?  
Time-reversal:  $f_{1T}^\perp|_{DY} = -f_{1T}^\perp|_{DIS}$        $h_1^\perp|_{DY} = -h_1^\perp|_{DIS}$   
(Collins, 2002)

# Universality of non-perturbative factors

## 1. Problem

- Semi-inclusive DIS

$$\sigma|_{DIS} \propto \hat{\sigma}_{pert} \otimes \text{pdf} \otimes \text{ff} \otimes \text{soft}$$

- Drell-Yan

$$\sigma|_{DY} \propto \hat{\sigma}_{pert} \otimes \text{pdf} \otimes \text{pdf} \otimes \text{soft}$$

$$\text{pdf}|_{DIS} \stackrel{?}{=} \text{pdf}|_{DY}$$

$$\text{soft}|_{DIS} \stackrel{?}{=} \text{soft}|_{DY}$$

## 2. Results

- Parton densities:

$$\text{Time-reversal: } f_{1T}^\perp|_{DY} = -f_{1T}^\perp|_{DIS} \qquad h_1^\perp|_{DY} = -h_1^\perp|_{DIS}$$

6 T-even pdfs are universal

- Soft factor:

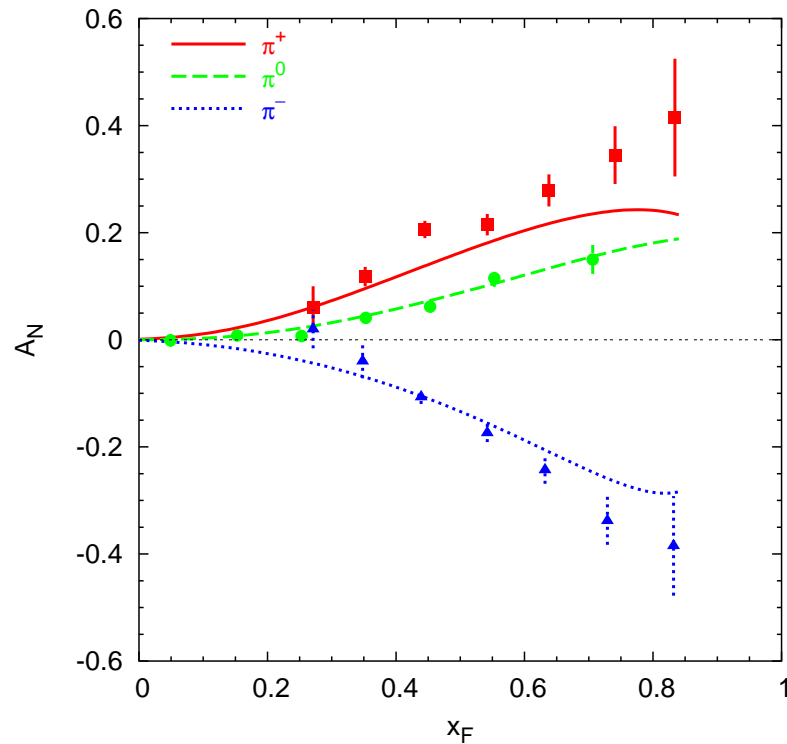
$$\text{Time-reversal and analytical structure: } \text{soft}|_{DIS} = \text{soft}|_{DY}$$

(Metz 2002; Collins, Metz, 2004)

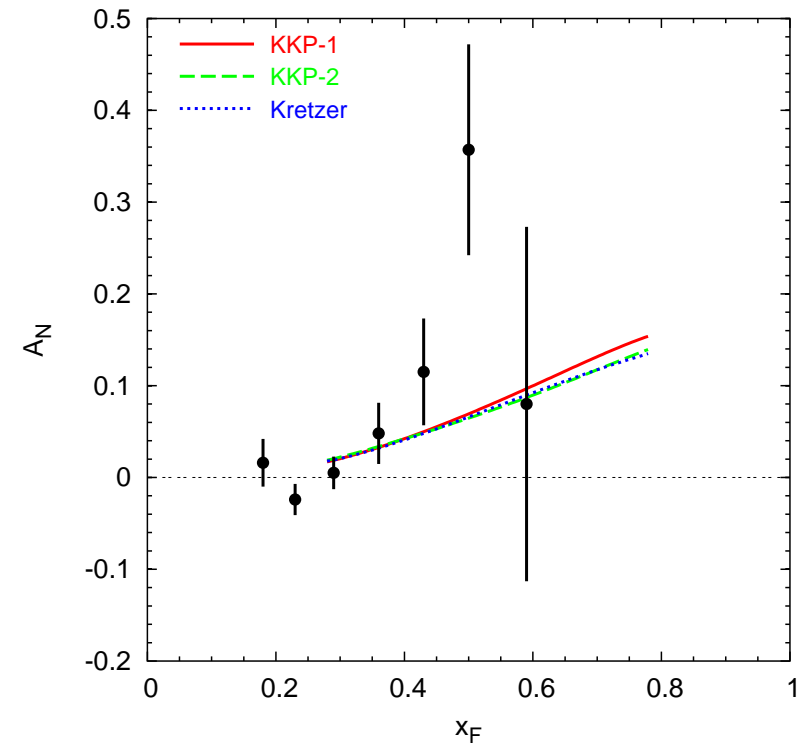
# Transverse SSA in $pp^\uparrow \rightarrow \pi X$

$$A_N = \frac{d\sigma^\uparrow - d\sigma^\downarrow}{d\sigma^\uparrow + d\sigma^\downarrow}$$

$$x_F = \frac{2P_{hL}}{\sqrt{s}}$$



FermiLab, E704, 1990  $\sqrt{s} = 20$  GeV



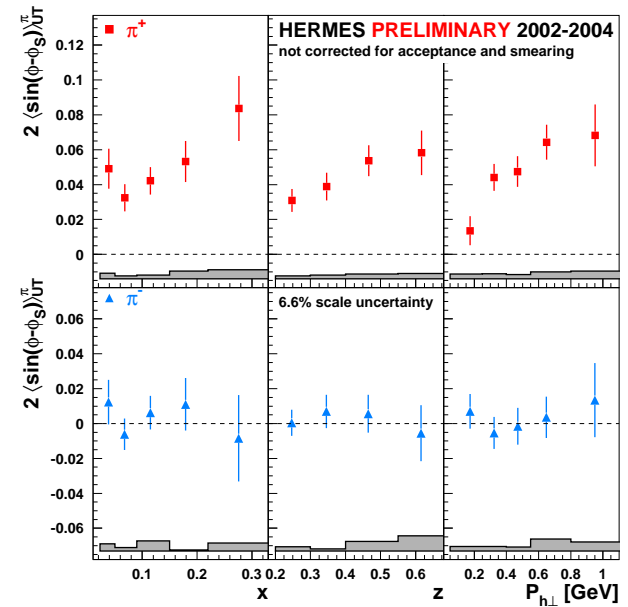
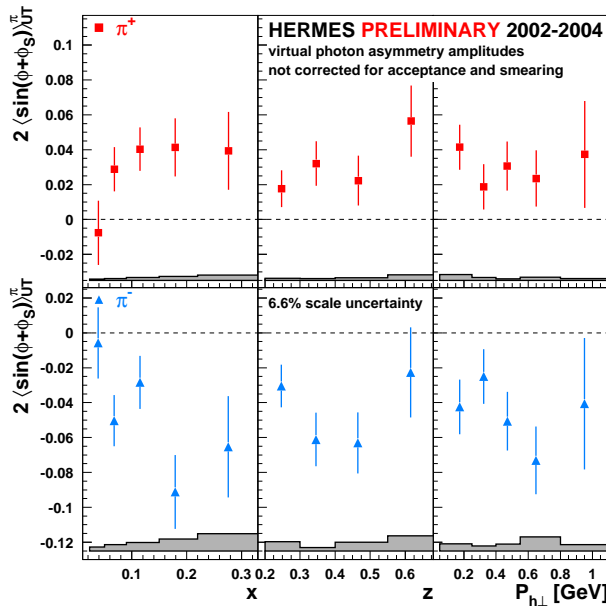
RHIC, STAR, 2004  $\sqrt{s} = 200$  GeV

SSA vanishes in leading twist collinear pQCD-description  
 $\rightarrow$  (effective)  $p_T$ -dependent formalism, etc.

# SSAs in semi-inclusive DIS

- Longitudinal SSAs (twist-3,  $\propto 1/Q$ )
  - Target spin asymmetry  $A_{UL}$ : data from HERMES
  - Beam spin asymmetry  $A_{LU}$ : data from CLAS
- Transverse SSA (leading twist)

$$A_{UT} \propto \sin(\Phi_h + \Phi_S) h_1(x) H_1^\perp(z) + \sin(\Phi_h - \Phi_S) f_{1T}^\perp(x) D_1(z) + \sin(3\Phi_h + \Phi_S) \dots + \mathcal{O}(1/Q)$$



Also COMPASS data on  $l D^\uparrow \rightarrow l' H X$

# Phenomenology of the Sivers effect

(Collins, Efremov, Goeke, Grosse-Perdekamp, Menzel, Meredith,  
Metz, Schweitzer, 2004, 2005)

Focus on Sivers effect in semi-inclusive DIS and Drell-Yan

Effect also in  $p^\uparrow \bar{p} \rightarrow H X$ ,  $p^\uparrow \bar{p} \rightarrow H_1 H_2 X$  etc. (theoretically more involved)

1. Extraction of Sivers function  $f_{1T}^\perp$  from data on  $A_{UT}$  ( $e p^\uparrow \rightarrow e \pi X$ )

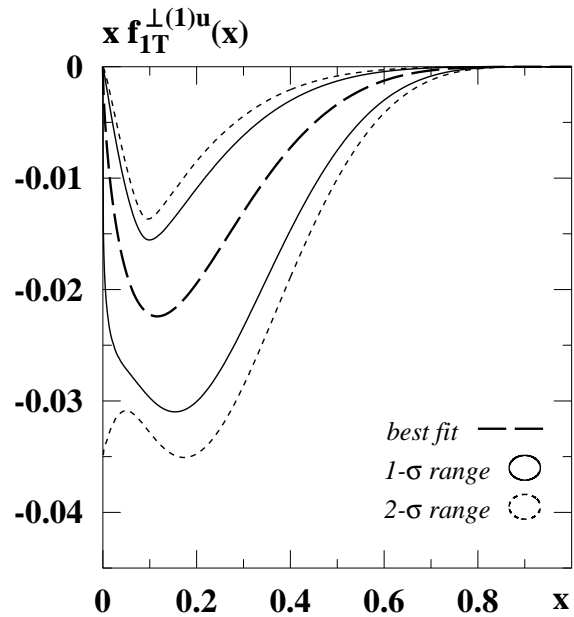
Input/approximations:

- HERMES data
- Neglect of soft gluon emission
- Gaussian  $p_T$ -behaviour, e.g.,

$$f_1(x, \vec{p}_T^2) = \frac{1}{\pi \langle p_T^2 \rangle} f_1(x) e^{-\vec{p}_T^2 / \langle p_T^2 \rangle}$$

- Exact result in large  $N_c$ -limit (Pobylitsa, 2003):

$$f_{1T}^{\perp u}(x, \vec{p}_T^2) = -f_{1T}^{\perp d}(x, \vec{p}_T^2)$$



$$f_{1T}^{\perp(1)}(x) \equiv \int d^2 \vec{p}_T \frac{\vec{p}_T^2}{2M^2} f_{1T}^{\perp}(x, \vec{p}_T^2)$$

### Extractions also by

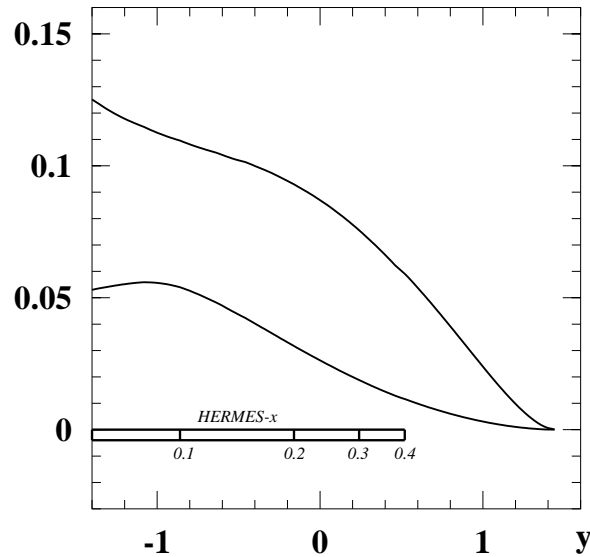
- Anselmino, Boglione, D' Alesio, Kotzinian, Murgia, Prokudin, 2005
- Vogelsang, Yuan, 2005

Comparison in: Anselmino et al., hep-ph/0511017

## 2. Calculation of the Sivers asymmetry for Drell-Yan

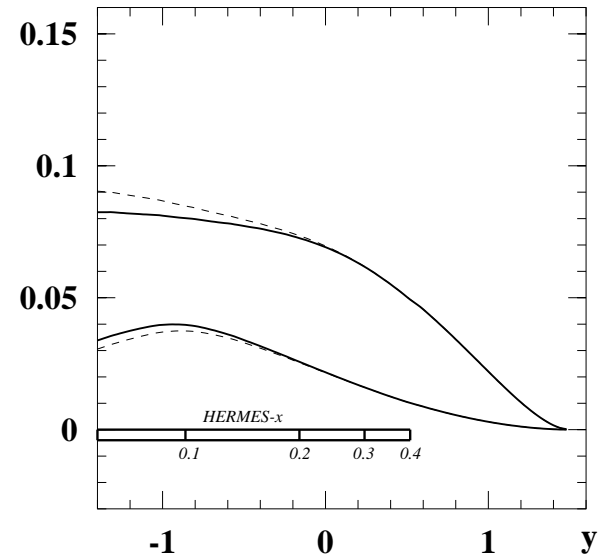
$$y = \frac{1}{2} \ln \frac{x_1}{x_2} \qquad f_{1T}^{\perp(1)\bar{q}}(x) = \pm \frac{f_1^{\bar{u}}(x) + f_1^{\bar{d}}(x)}{f_1^u(x) + f_1^d(x)} f_{1T}^{\perp(1)q}(x)$$

$A_{UT}^{\sin(\phi - \phi_S)}$  in  $p^\uparrow \bar{p} \rightarrow l^+ l^- X$  at PAX



GSI  $p^\uparrow \bar{p} \rightarrow l^+ l^- X$   
 $s = 45 \text{ GeV}^2 \quad Q^2 = 6.25 \text{ GeV}^2$

$A_{UT}^{\sin(\phi - \phi_S)}$  in  $p^\uparrow \pi^- \rightarrow l^+ l^- X$  at COMPASS

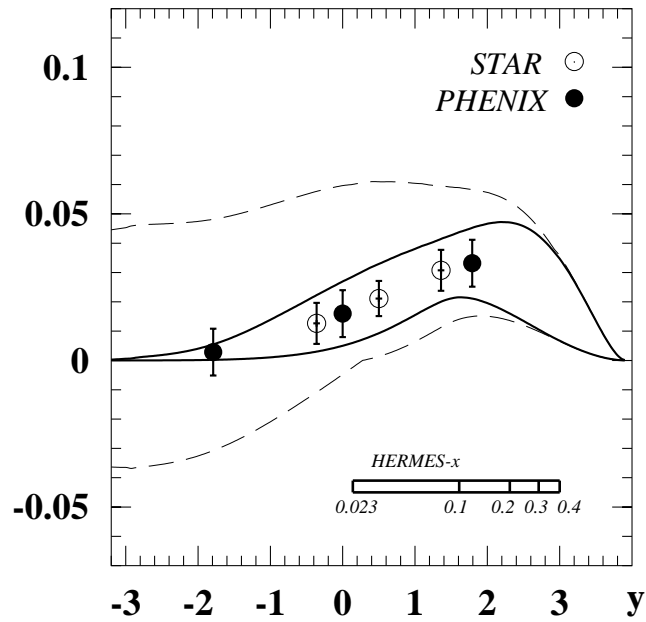


COMPASS  $p^\uparrow \pi^- \rightarrow l^+ l^- X$   
 $s = 400 \text{ GeV}^2 \quad Q^2 = 20 \text{ GeV}^2$

→ Prediction  $f_{1T}^{\perp}|_{DY} = -f_{1T}^{\perp}|_{DIS}$  can be checked experimentally

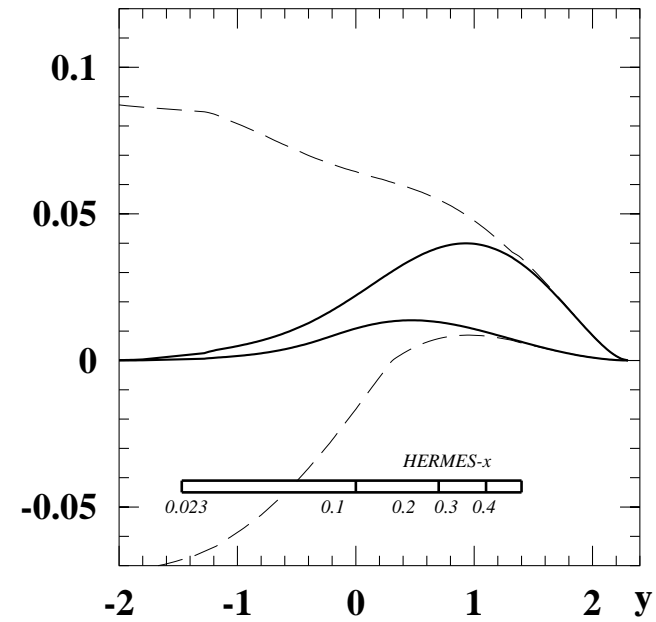


$A_{UT}^{\sin(\phi - \phi_S)}$  in  $p^\uparrow p \rightarrow l^+ l^- X$  at RHIC  $Q=4\text{GeV}$



RHIC  $p^\uparrow p \rightarrow l^+ l^- X$   
 $\sqrt{s} = 200 \text{ GeV}$   $Q^2 = 16 \text{ GeV}^2$

$A_{UT}^{\sin(\phi - \phi_S)}$  in  $p^\uparrow p \rightarrow l^+ l^- X$  at RHIC  $Q=20\text{GeV}$



RHIC  $p^\uparrow p \rightarrow l^+ l^- X$   
 $\sqrt{s} = 200 \text{ GeV}$   $Q^2 = 400 \text{ GeV}^2$

→ In  $pp$ -collisions strong sensitivity to  $f_{1T}^{\perp \bar{q}}$

# Drell-Yan measurements in $p\bar{p}$ -interactions

## 1. Double-polarized

- Transverse polarization:

$$A_{TT}^{p\bar{p}} = a_{TT} \frac{\sum_q e_q^2 (h_1^{q/p}(x_1, Q^2) h_1^{\bar{q}/\bar{p}}(x_2, Q^2)) + (x_1 \leftrightarrow x_2)}{\sum_q e_q^2 (f_1^{q/p}(x_1, Q^2) f_1^{\bar{q}/\bar{p}}(x_2, Q^2)) + (x_1 \leftrightarrow x_2)}$$

→ Unique tool to measure transversity  $h_1$

- Longitudinal polarization:  $A_{LL}^{p\bar{p}} \propto g_1^{q/p} g_1^{\bar{q}/\bar{p}}$

## 2. Single-polarized

- Transverse polarization:

$$A_{TU} \propto \sin(\Phi_q - \Phi_S) f_{1T}^\perp(x_1) f_1(x_2) + \sin(\Phi_q + \Phi_S) h_1(x_1) h_1^\perp(x_2) + \dots$$

→ Access to  $f_{1T}^\perp$  and  $h_1^\perp$

- Longitudinal polarization:  $A_{LU} \propto \sin(2\Phi_q) h_{1L}^\perp(x_1) h_1^\perp(x_2)$

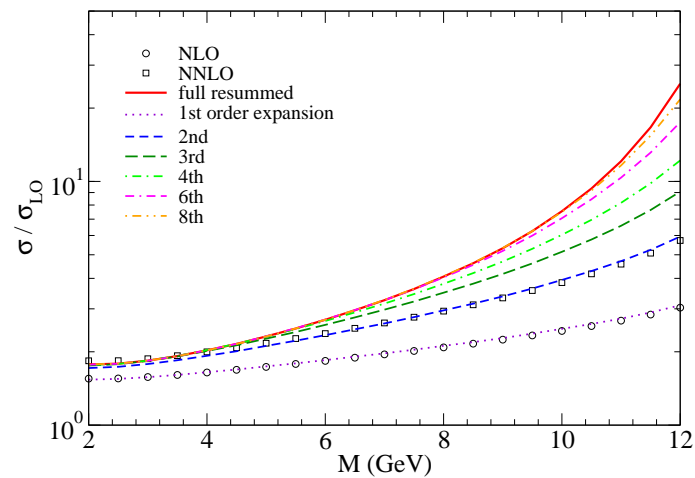
### 3. Unpolarized

- Unpolarized cross section: fixed order vs. resummation

$$z = \frac{Q^2}{x_1 x_2 s} = \frac{Q^2}{\hat{s}} \quad \left\{ \begin{array}{l} = 1 \quad \text{at tree level} \\ < 1 \quad \text{at } \mathcal{O}(\alpha_s) \text{ and beyond} \end{array} \right.$$

$$k\text{-loop calculation: } \alpha_s^k \left( \frac{\ln^{2k-1}(1-z)}{1-z} \right)_+$$

$\sqrt{s} = 14.5 \text{ GeV}$  (Shimizu, Sterman, Vogelsang, Yokaya, 2005)



- Angular distribution

$$\frac{d\sigma}{d\Omega} \propto 1 + \lambda \cos^2 \Theta + \mu \sin 2\Theta \cos \Phi + \frac{\nu}{2} \sin^2 \Theta \cos 2\Phi$$

Collinear picture at  $\mathcal{O}(\alpha_s)$ :  $\lambda + 2\nu = 1$

(Lam, Tung, 1978)

- (1) hardly modified at  $\mathcal{O}(\alpha_s^2)$
- (2) strongly violated by data,  $\nu$  large

Possible explanation:

(Boer, 1999; Boer, Brodsky, Hwang, 2002; ...)

$$\nu \propto h_1^\perp(x_1) h_1^\perp(x_2)$$

Relation of T-odd pdfs to generalized parton distributions:

(Burkardt, 2004, 2005, 2006)

$$f_{1T}^\perp \propto E \qquad h_1^\perp \propto 2H_T + E_T$$

## Summary

- Siverson function can be non-zero
- Universality of non-perturbative factors
  1.  $f_{1T}^\perp|_{DY} = -f_{1T}^\perp|_{DIS}$        $h_1^\perp|_{DY} = -h_1^\perp|_{DIS}$
  2.  $\text{soft}|_{DIS} = \text{soft}|_{DY}$
- Sign-reversal of Siverson function can be checked in  $p^\uparrow \bar{p} \rightarrow l^+ l^- X$  and/or  $p \bar{p}^\uparrow \rightarrow l^+ l^- X$
- Drell-Yan measurements (DY is cleanest hard  $HH$ -collision process)
  1. Information on nucleon structure:  $h_1, g_1, f_{1T}^\perp, h_1^\perp, \dots$
  2. Application of pQCD to  $p \bar{p} \rightarrow l^+ l^- X, HH \rightarrow \dots$

Dedicated Drell-Yan program for proton-antiproton interactions would be valuable!